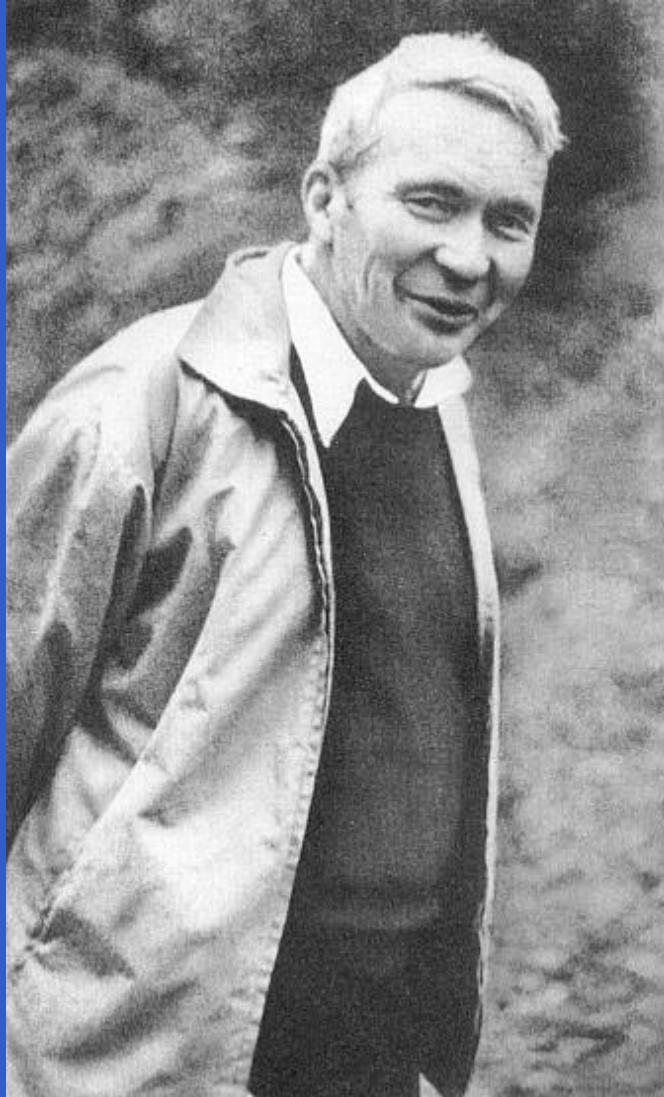


Andrei Nikolaevich Kolmogorov



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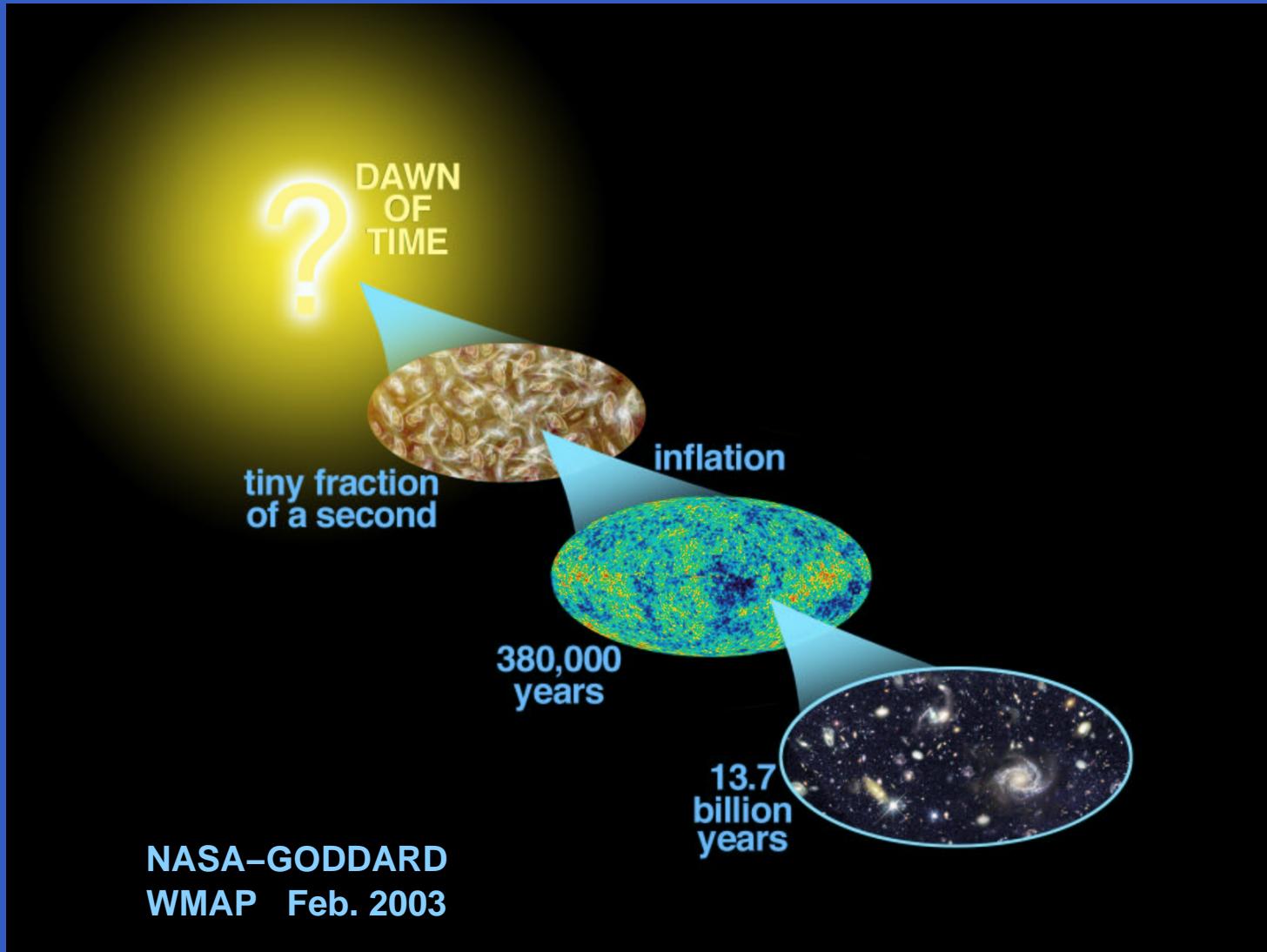
Back to the primordial universe by a Monge–Ampère–Kantorovich mass transportation method

Uriel FRISCH

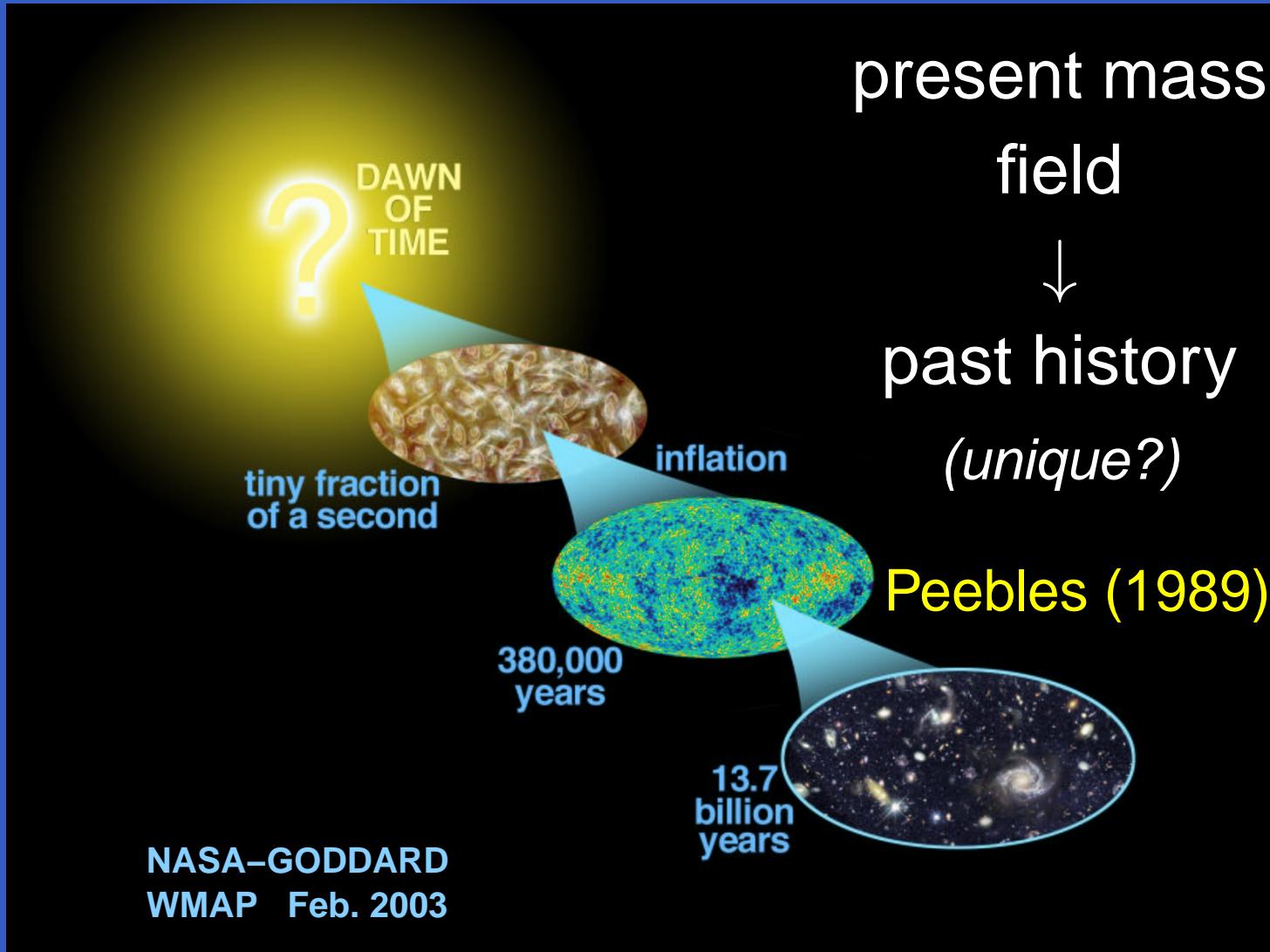
Observatoire de la Côte d'Azur, Nice, France

- Y. Brenier, U. Frisch, M. Hénon, G. Loeper, S. Matarrese, R. Mohayaee,
A. Sobolevskiĭ *Mon. Not. R. Astron. Soc.* (2003, submitted), **astro-ph/0304214**
- U. Frisch, S. Matarrese, R. Mohayaee, A. Sobolevski *Nature* **417** (2002) 260–262

Brief history of Universe



History of Universe reconstructed



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Fluid dynamics in expanding universe

Euler:

$$\partial_\tau \mathbf{v} + (\mathbf{v} \cdot \nabla_{\mathbf{x}}) \mathbf{v} = -\frac{3}{2\tau} (\mathbf{v} + \nabla_{\mathbf{x}} \varphi_g)$$

Mass conservation:

$$\partial_\tau \rho + \nabla_{\mathbf{x}} \cdot (\rho \mathbf{v}) = 0$$

Poisson:

$$\nabla_{\mathbf{x}}^2 \varphi_g = \frac{\rho - 1}{\tau}$$

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Fluid dynamics in expanding universe

Euler:

$$\partial_\tau \mathbf{v} + (\mathbf{v} \cdot \nabla_{\mathbf{x}}) \mathbf{v} = -\frac{3}{2\tau} (\mathbf{v} + \nabla_{\mathbf{x}} \varphi_g)$$

Mass conservation:

$$\partial_\tau \rho + \nabla_{\mathbf{x}} \cdot (\rho \mathbf{v}) = 0$$

Poisson:

$$\nabla_{\mathbf{x}}^2 \varphi_g = \frac{\rho - 1}{\tau}$$

“Slaving” as $\tau \rightarrow 0$:

$$\mathbf{v}_{\text{in}} = -\nabla_{\mathbf{x}} \varphi_g^{(\text{in})}, \quad \rho_{\text{in}} = 1$$

Fluid dynamics in expanding universe

Euler:

$$\partial_\tau \mathbf{v} + (\mathbf{v} \cdot \nabla_{\mathbf{x}}) \mathbf{v} = -\frac{3}{2\tau} (\mathbf{v} + \nabla_{\mathbf{x}} \varphi_g)$$

Mass conservation:

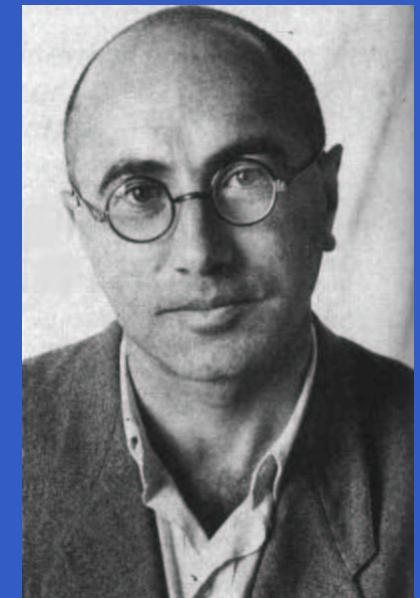
$$\partial_\tau \rho + \nabla_{\mathbf{x}} \cdot (\rho \mathbf{v}) = 0$$

Poisson:

$$\nabla_{\mathbf{x}}^2 \varphi_g = \frac{\rho - 1}{\tau}$$

Zeldovich approximation:

$$\mathbf{v} \equiv -\nabla_{\mathbf{x}} \varphi_g$$



late 1960s

Potential Lagrangian map

“Zeldovich” equation:

$$\partial_\tau \mathbf{v} + \mathbf{v} \cdot \nabla_{\mathbf{x}} \mathbf{v} = 0$$

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Potential Lagrangian map

“Zeldovich” equation:

$$\partial_\tau \mathbf{v} + \mathbf{v} \cdot \nabla_{\mathbf{x}} \mathbf{v} = 0$$

$$\mathbf{v}_{\text{in}} = -\nabla_{\mathbf{x}} \varphi_g^{(\text{in})} \Rightarrow \mathbf{v} = -\nabla_{\mathbf{x}} \varphi_v$$

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Potential Lagrangian map

“Zeldovich”/Burgers equation:

$$\partial_\tau \mathbf{v} + \mathbf{v} \cdot \nabla_{\mathbf{x}} \mathbf{v} = 0$$

$$\mathbf{v}_{\text{in}} = -\nabla_{\mathbf{x}} \varphi_g^{(\text{in})} \Rightarrow \mathbf{v} = -\nabla_{\mathbf{x}} \varphi_v$$

Potential Lagrangian map

“Zeldovich”/Burgers equation:

$$\partial_\tau \mathbf{v} + \mathbf{v} \cdot \nabla_{\mathbf{x}} \mathbf{v} = 0$$

$$\mathbf{v}_{\text{in}} = -\nabla_{\mathbf{x}} \varphi_g^{(\text{in})} \Rightarrow \mathbf{v} = -\nabla_{\mathbf{x}} \varphi_v$$

$$\mathbf{x} = \mathbf{q} + \tau \mathbf{v}_{\text{in}}(\mathbf{q}) = \mathbf{q} - \tau \nabla \varphi_g^{(\text{in})}(\mathbf{q})$$

$$= \nabla \left[\frac{|\mathbf{q}|^2}{2} - \tau \varphi_g^{(\text{in})}(\mathbf{q}) \right]$$

$$= \nabla \Phi(\mathbf{q})$$

Bertschinger–Dekel (1989)

Potential Lagrangian map

“Zeldovich”/Burgers equation:

$$\partial_\tau \mathbf{v} + \mathbf{v} \cdot \nabla_{\mathbf{x}} \mathbf{v} = 0$$

$$\mathbf{v}_{\text{in}} = -\nabla_{\mathbf{x}} \varphi_g^{(\text{in})} \Rightarrow \mathbf{v} = -\nabla_{\mathbf{x}} \varphi_v$$

$$\mathbf{x} = \mathbf{q} + \tau \mathbf{v}_{\text{in}}(\mathbf{q}) = \mathbf{q} - \tau \nabla \varphi_g^{(\text{in})}(\mathbf{q})$$

$$= \nabla \left[\frac{|\mathbf{q}|^2}{2} - \tau \varphi_g^{(\text{in})}(\mathbf{q}) \right]$$

$$= \nabla \Phi(\mathbf{q}) \quad \dots \text{(graph) invertible if } \Phi \text{ is convex}$$

Bertschinger–Dekel (1989)

The Monge–Ampère equation

Mass conservation:

$$\det \nabla_{\mathbf{q}} \mathbf{x} = \frac{\rho_{\text{in}}(\mathbf{q})}{\rho_0(\mathbf{x}(\mathbf{q}))} = \frac{\text{initial density}}{\text{present density}}$$

The Monge–Ampère equation

Mass conservation, $\rho_{\text{in}} = 1$:

$$\det \nabla_{\mathbf{q}} \mathbf{x} = \frac{1}{\rho_0(\mathbf{x}(\mathbf{q}))} = \frac{\text{initial density}}{\text{present density}}$$

The Monge–Ampère equation

Mass conservation, $\rho_{\text{in}} = 1$:

$$\det \nabla_{\mathbf{q}} \mathbf{x} = \frac{1}{\rho_0(\mathbf{x}(\mathbf{q}))} = \frac{\text{initial density}}{\text{present density}}$$

As $\mathbf{x}(\mathbf{q}) = \nabla_{\mathbf{q}} \Phi(\mathbf{q})$:

$$\det(\nabla_{q_i} \nabla_{q_j} \Phi(\mathbf{q})) = \frac{1}{\rho_0(\nabla_{\mathbf{q}} \Phi(\mathbf{q}))}$$

The Monge–Ampère equation

Mass conservation, $\rho_{\text{in}} = 1$:

$$\det \nabla_{\mathbf{q}} \mathbf{x} = \frac{1}{\rho_0(\mathbf{x}(\mathbf{q}))} = \frac{\text{initial density}}{\text{present density}}$$

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Legendre–Fenchel transform:

$$\Theta(\mathbf{x}) = \max_{\mathbf{q}} (\mathbf{x} \cdot \mathbf{q} - \Phi(\mathbf{q}))$$

The Monge–Ampère equation

Mass conservation, $\rho_{\text{in}} = 1$:

$$\det \nabla_{\mathbf{x}} \mathbf{q} = \frac{\rho_0(\mathbf{x})}{1}$$

As $\mathbf{q}(\mathbf{x}) = \nabla_{\mathbf{x}} \Theta(\mathbf{x})$:

$$\det(\nabla_{x_i} \nabla_{x_j} \Theta(\mathbf{x})) = \rho_0(\mathbf{x})$$



1820

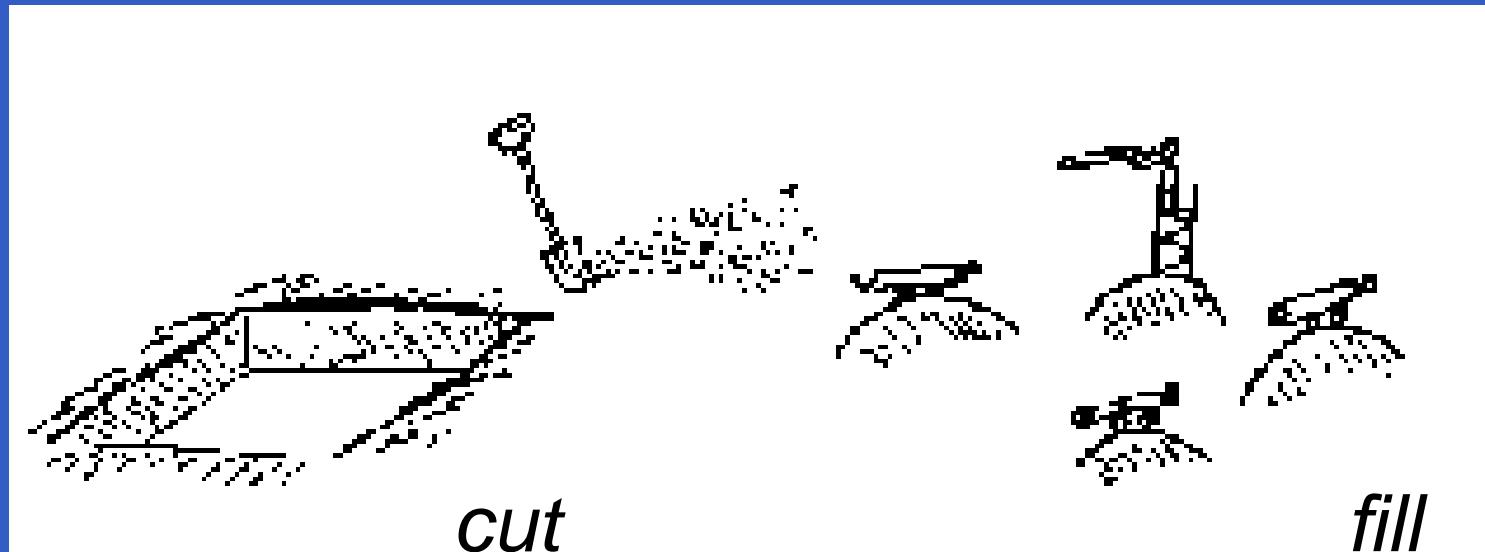
$\rho_0(\mathbf{x})$ prescribed

Monge's mass transportation problem



*...Il n'est pas indifférent que telle molécule de déblai soit transportée dans tel ou tel autre endroit du remblai, mais qu'il y a une certaine distribution à faire des molécules du premier dans le second, d'après laquelle la somme de ces produits sera la moindre possible, et le prix du transport total sera un **minimum**.*

Monge's mass transportation problem



*...It is not indifferent that any given molecule of the cuts be transported to this or that place in the fills, but there ought to be a certain distribution of molecules of the former into the latter, according to which the sum of these products will be the least possible, and the cost of transportation will be a **minimum**.*

Monge's mass transportation problem



For given $\rho_{\text{in}}(q)$, $\rho_0(x)$ minimize

$$\int |\mathbf{x}(q) - q| \rho_{\text{in}}(q) dq = \int |\mathbf{x} - q(\mathbf{x})| \rho_0(\mathbf{x}) d\mathbf{x}$$

over all $(\mathbf{x}(q), q(\mathbf{x}))$ such that $\rho_{\text{in}}(q) dq = \rho_0(\mathbf{x}) d\mathbf{x}$

Monge, Ampère, mass transportation

For given $\rho_{\text{in}}(\mathbf{q})$, $\rho_0(\mathbf{x})$ minimize

$$\int |\mathbf{x}(\mathbf{q}) - \mathbf{q}|^2 \rho_{\text{in}}(\mathbf{q}) d\mathbf{q} = \int |\mathbf{x} - \mathbf{q}(\mathbf{x})|^2 \rho_0(\mathbf{x}) d\mathbf{x}$$

over all $(\mathbf{x}(\mathbf{q}), \mathbf{q}(\mathbf{x}))$ such that $\rho_{\text{in}}(\mathbf{q}) d\mathbf{q} = \rho_0(\mathbf{x}) d\mathbf{x}$

Monge, Ampère, mass transportation

Theorem (Brenier 1987, 1991) *The minimizing maps are gradients of convex functions:*

$$\mathbf{x}(\mathbf{q}) = \nabla_{\mathbf{q}} \Phi(\mathbf{q}), \quad \mathbf{q}(\mathbf{x}) = \nabla_{\mathbf{x}} \Theta(\mathbf{x})$$

Φ and Θ solve suitable Monge–Ampère equations

For given $\rho_{\text{in}}(\mathbf{q})$, $\rho_0(\mathbf{x})$ minimize

$$\int |\mathbf{x}(\mathbf{q}) - \mathbf{q}|^2 \rho_{\text{in}}(\mathbf{q}) d\mathbf{q} = \int |\mathbf{x} - \mathbf{q}(\mathbf{x})|^2 \rho_0(\mathbf{x}) d\mathbf{x}$$

over all $(\mathbf{x}(\mathbf{q}), \mathbf{q}(\mathbf{x}))$ such that $\rho_{\text{in}}(\mathbf{q}) d\mathbf{q} = \rho_0(\mathbf{x}) d\mathbf{x}$

Kantorovich relaxation

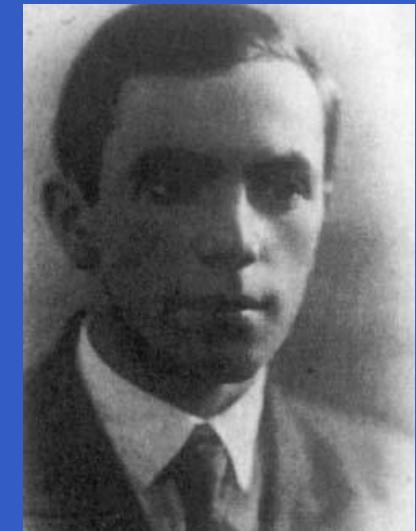
For given $\rho_{\text{in}}(\mathbf{q})$, $\rho_0(\mathbf{x})$ minimize

$$\int |\mathbf{x} - \mathbf{q}|^2 \rho(\mathbf{q}, \mathbf{x}) d\mathbf{q} d\mathbf{x}$$

over all $\rho(\mathbf{q}, \mathbf{x})$ such that

$$\int \rho(\mathbf{q}, \mathbf{x}) d\mathbf{x} = \rho_{\text{in}}(\mathbf{q})$$

$$\int \rho(\mathbf{q}, \mathbf{x}) d\mathbf{q} = \rho_0(\mathbf{x})$$



1942

Discretization and assignment

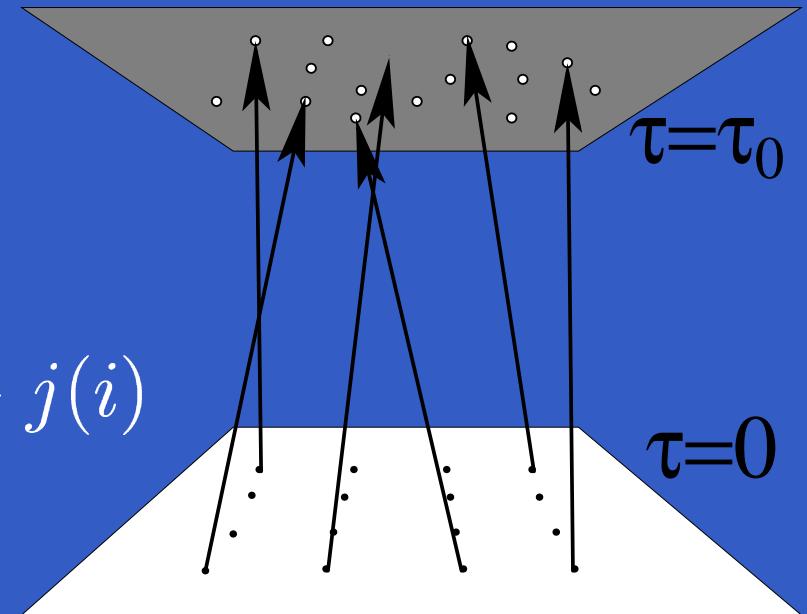
Discrete densities:

$$\rho_0(\mathbf{x}) = \sum_{i=1}^N \delta(\mathbf{x} - \mathbf{x}_i), \quad \rho_{\text{in}}(\mathbf{q}) = \sum_{j=1}^N \delta(\mathbf{q} - \mathbf{q}_j)$$

Minimize the cost

$$\sum_{i=1}^N |\mathbf{x}_i - \mathbf{q}_{j(i)}|^2$$

over all permutations $i \mapsto j(i)$
of $\{1, 2, \dots, N\}$



Relaxation and the dual problem

$$\text{Minimize} \sum_{i=1}^N |\mathbf{x}_i - \mathbf{q}_{j(i)}|^2$$

over all permutations $i \mapsto j(i)$

Relaxation and the dual problem

$$\text{Minimize} \sum_{i,j=1}^N |\mathbf{x}_i - \mathbf{q}_j|^2 f_{ij}$$

over all bistochastic matrices:

$$f_{ij} \geq 0, \quad \sum_{k=1}^N f_{kj} = \sum_{k=1}^N f_{ik} = 1$$

Relaxation and the dual problem

$$\text{Minimize} \sum_{i,j=1}^N |\mathbf{x}_i - \mathbf{q}_j|^2 f_{ij}$$

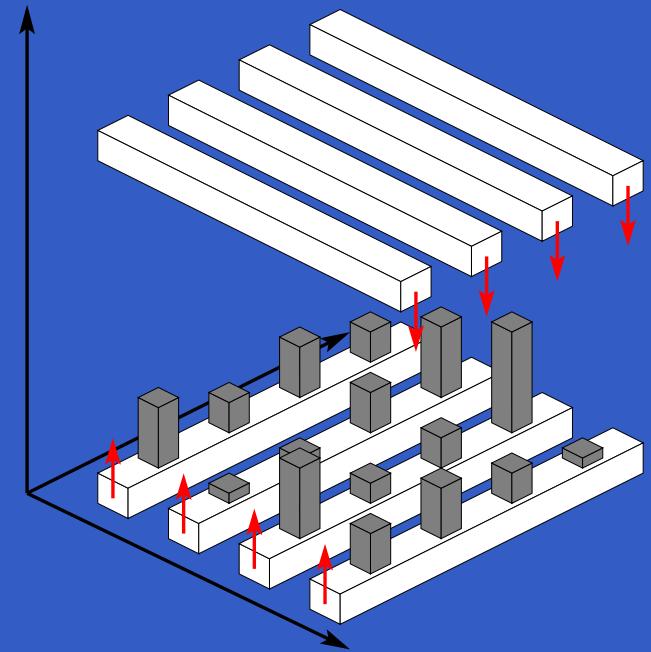
over all bistochastic matrices:

$$f_{ij} \geq 0, \quad \sum_{k=1}^N f_{kj} = \sum_{k=1}^N f_{ik} = 1$$

Dual problem:

$$\text{Minimize} \sum_{i=1}^N \alpha_i - \sum_{j=1}^N \beta_j$$

$$\alpha_i - \beta_j \geq C - |\mathbf{x}_i - \mathbf{q}_j|^2$$



Hénon 1950s–1990s

Relaxation and the dual problem

$$\text{Minimize} \sum_{i,j=1}^N |\mathbf{x}_i - \mathbf{q}_j|^2 f_{ij}$$

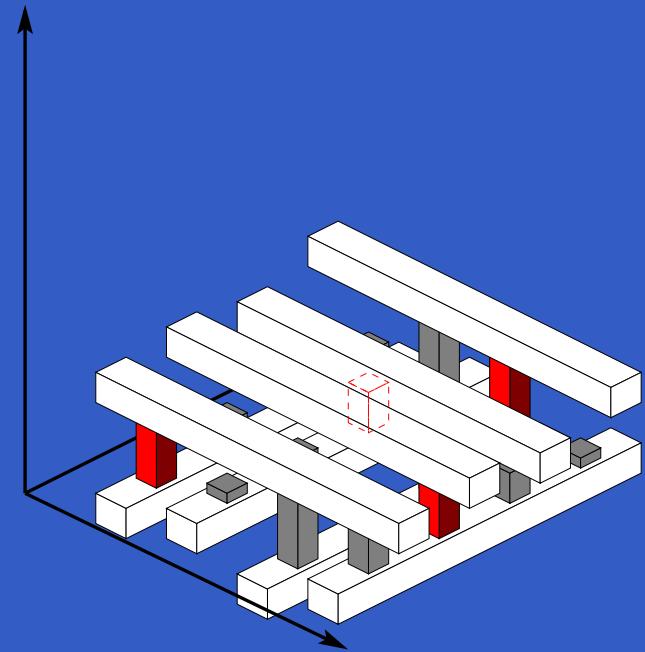
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Hénon 1950s–1990s

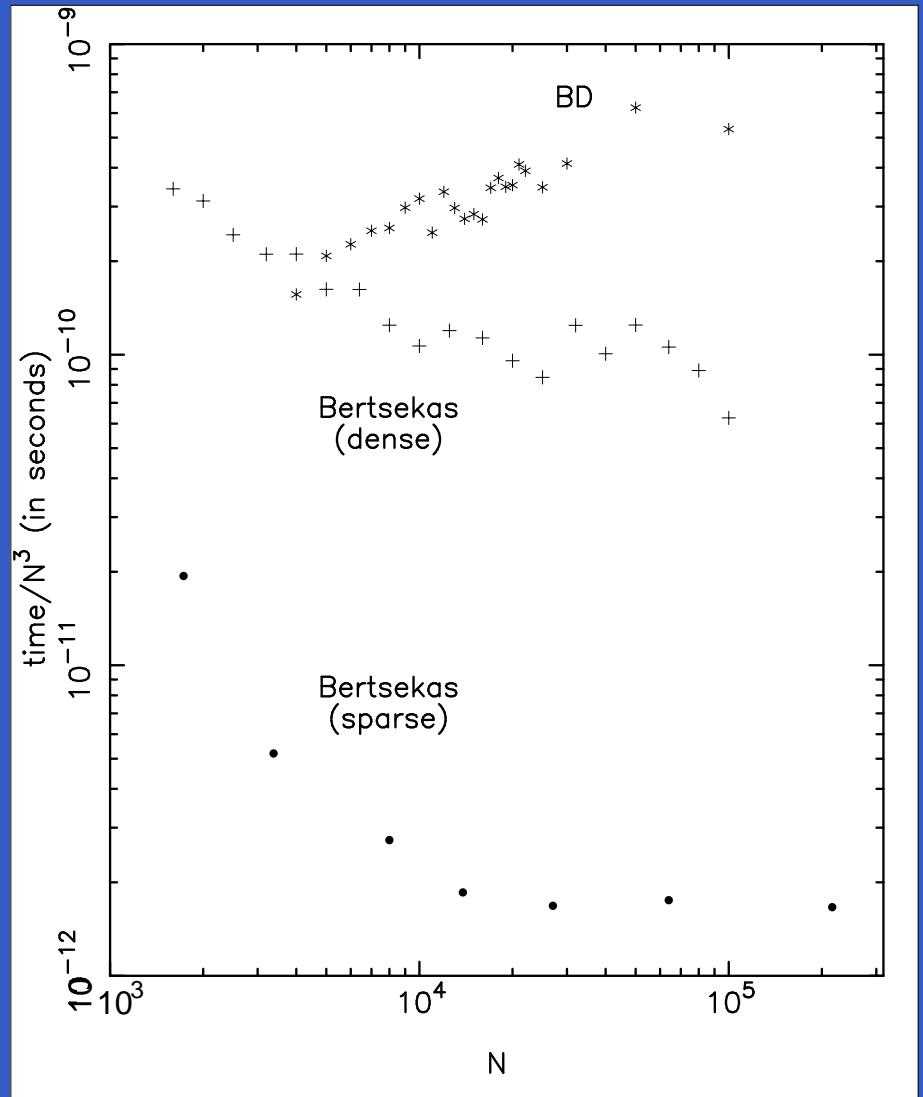
Time complexity

* Burkard & Derigs 1980

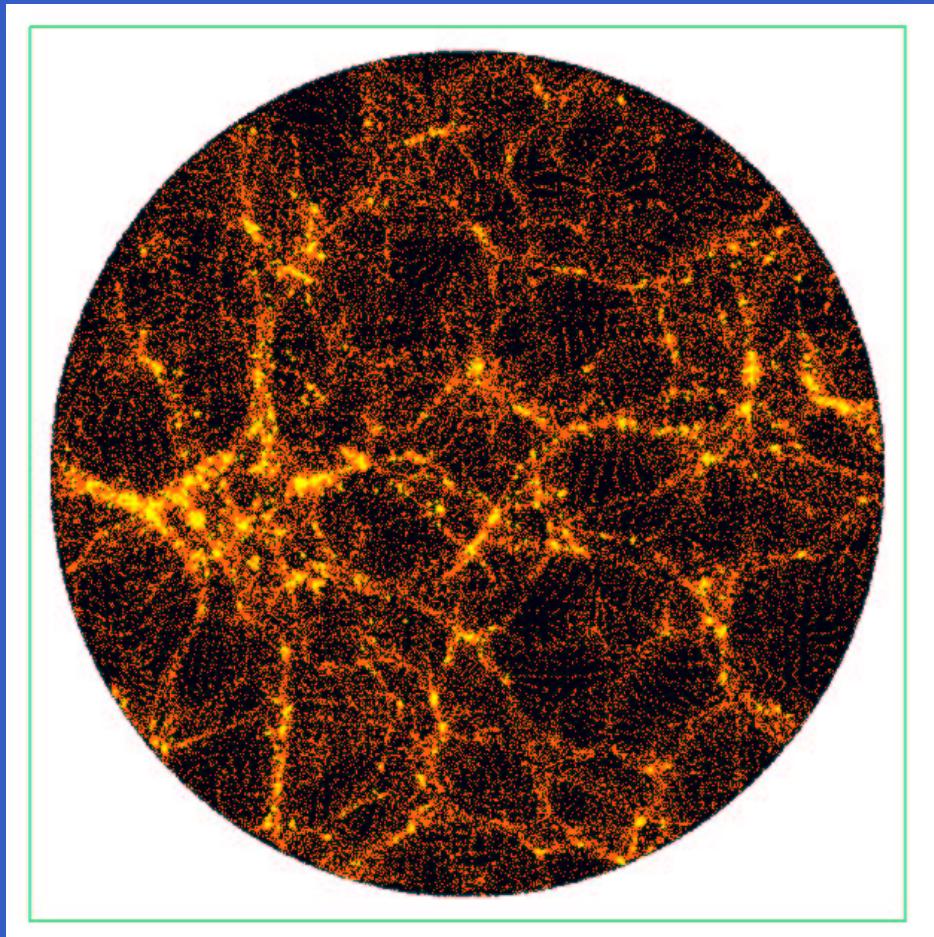
Bertsekas 1979–2003

- + “dense” auction
- “sparse” auction

(time in seconds
divided by N^3)

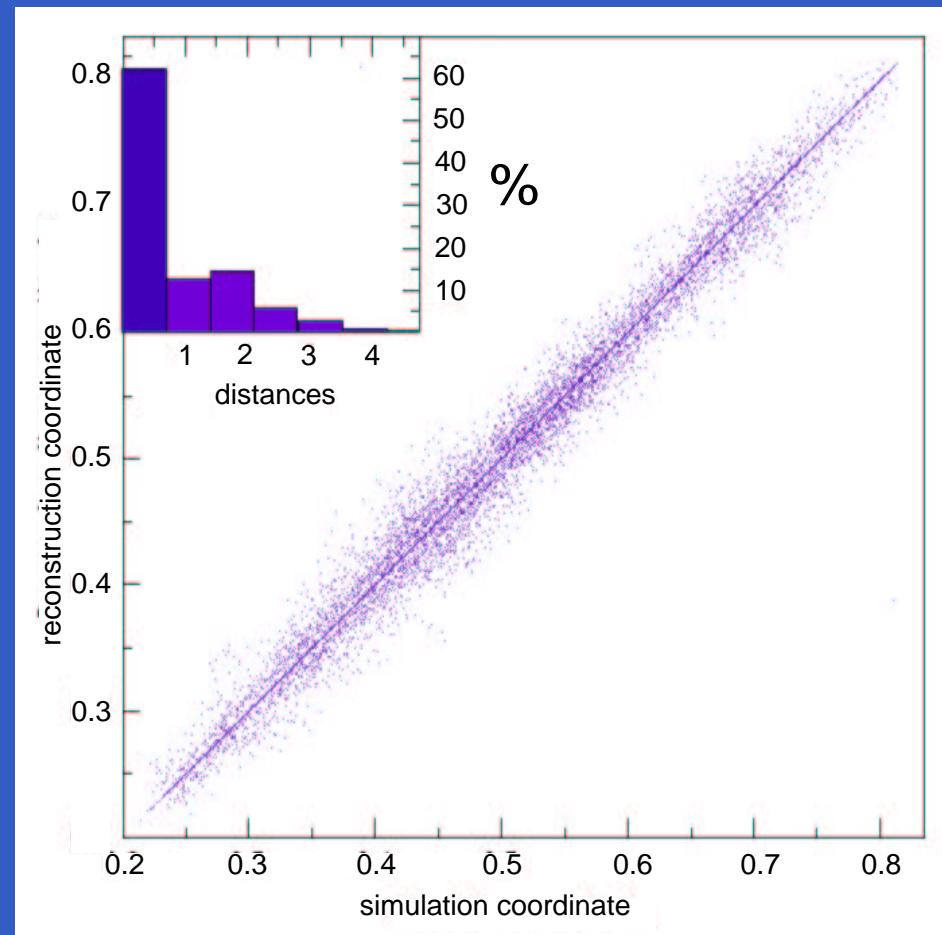
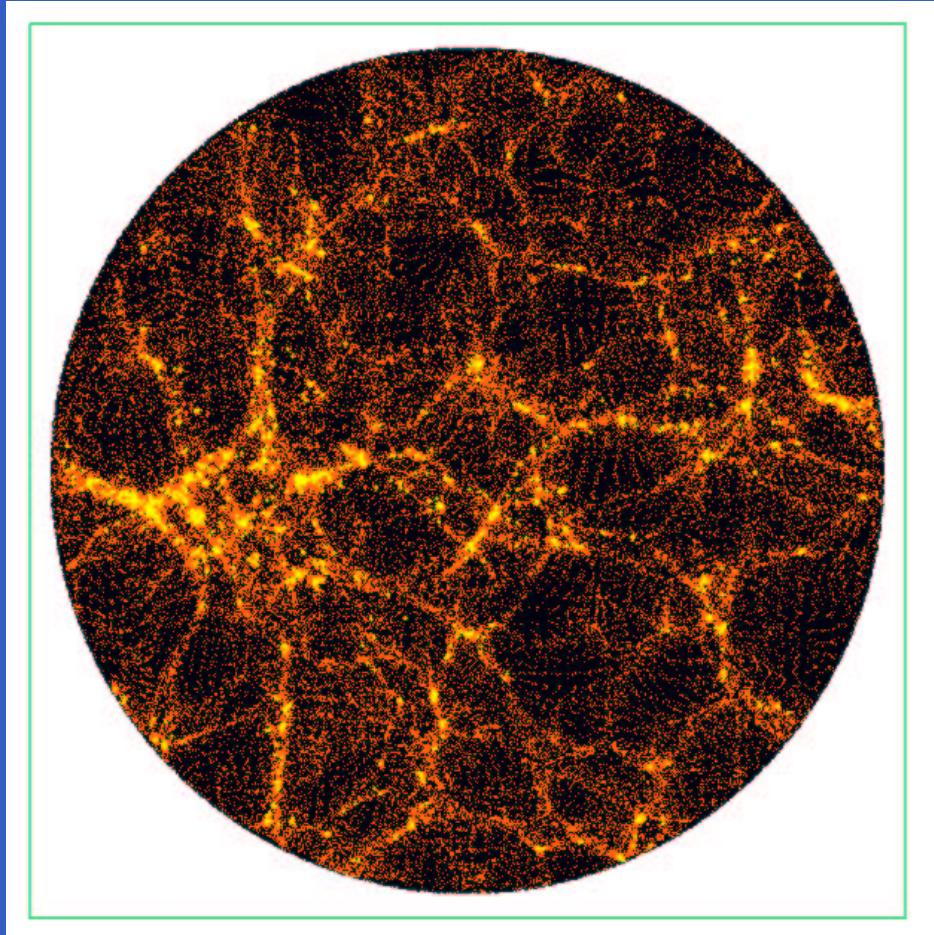


Testing against N -body simulation



128^3 points in a box
of $200 h^{-1}$ Mpc size
(middle 10% slice)

Testing against N -body simulation



Euler–Poisson variational problem

$$\text{Minimize} \int_0^{\tau_0} d\tau \int d^3x \tau^{3/2} \left(\rho |\mathbf{v}|^2 + \frac{3}{2} |\nabla_{\mathbf{x}} \varphi_g|^2 \right)$$

- Mass conservation: $\partial_\tau \rho + \nabla_{\mathbf{x}} \cdot (\rho \mathbf{v}) = 0$
- Poisson equation: $\nabla_{\mathbf{x}}^2 \varphi_g = \frac{\rho - 1}{\tau}$
- $\rho(\mathbf{x}, 0) = \rho_{\text{in}}(\mathbf{x}), \rho(\mathbf{x}, \tau_0) = \rho_0(\mathbf{x})$

Theorem (Loeper 2003) *Up to a change of variables this is a convex minimization problem with a unique solution $(\rho, \mathbf{v}, \varphi_g)$.*