Global models of planetary system formation

Richard Nelson Queen Mary, University of London









Occurence rates

30-50% of FGK stars host an Earth, super-Earth or Neptune with orbital period < 100 days (Fressin+ 2015) Microlensing survey results -> every star hosts \geq 1 Neptune-mass planet beyond the snowlike (Sumi et al 2016)



Eccentricities of Kepler super-Earths/Neptunes: e ~ 0.01 for 75% of planets e ~ 0.1-0.4 for 25% of planets (Wu & Lithwick 2013)

Kepler systems





Kepler systems



>6

Gas Glant

Kepler systems





Kepler 11



Dynamically stable

- but small perturbations away from observed planet orbital locations render system unstable (Mahajan & Wu 2014)



Orbital spacings of multiple systems



Period ratios of multiple systems











Kepler 223:

4 planets in resonance with period ratio 8:6:4:3 (Mills et al 2016)

Also Trappist-1: 7 planets with period ratios 8:5, 5:3, 3:2, 3:2, 4:3, 3:2





Radial velocity surveys also reveal multi planet systems with diverse properties.



Wasp 47 and similar systems



Planet	Radius	Mass	а	Period
WASP-47e	1.829	< 22	0.02	0.79
WASP-47b	12.77	~341 (286-414)	0.05	4.16
WASP-47d	3.63	15.2 ± 7	0.08	9.03
WASP-47c	-	394 ±70	1.36	572

Planet formation models - an unbiased perspective...

In situ formation: Migrate large amount of solids (100 M_{Earth}) into inner disc regions and grow via giant impacts

Pros: By construction can obtain systems with observed numbers of planets and their spacings

Cons: Formation will be rapid (< 1 Myr), but migration and influence of gas disc is ignored. Cannot explain single short-period planets.

Inside-out formation: Collect solids at disc inner edge and grow a sequence of planets one after the other

Pros: Can lead to systems of compact super-Earths as observed

Cons: Does not explain longer period systems -> treats short-period super-Earths as "special". MHD turbulence in inner disc regions maybe disruptive of planet formation.

Concurrent migration and growth: Build planets at range of distances and migrate when sufficiently massive

Pros: Can explain planetary systems with broad range of semi-major axes. Explains known resonant systems.

Cons: Generates too many resonant systems - require these to be unstable in the long-term.

Type I migration of low mass planets

Lindblad torque

- Gravitational interaction between planet and disc leads to the excitation of spiral density waves at Lindblad resonances (Goldreich & Tremaine 1978, 1980; Lin & Papaloizou 1979, 1984)
- Spiral wave exerts gravitational force on planet - removes angular momentum and drives inward migration
- Total Lindblad torque scales as:

 $\Gamma_0 = (q/h)^2 \Sigma_p r_p^4 \Omega_p^2$

- q = M_p / M_{*} planet/star mass ratio
- Σ = gas surface density
- Ω = angular velocity
- h = aspect ratio H/r
- $X_p = X$ at the planet location.

Σ(MMSN @1AU) => migration time [years] ≈ 1/q

→ 300000 yrs for an Earth , → 20000 yrs for a Neptune !



Lindblad torque

- Gravitational interaction between planet and disc leads to the excitation of spiral density waves at Lindblad resonances (Goldreich & Tremaine 1978, 1980; Lin & Papaloizou 1979, 1984)
- Spiral wave exerts gravitational force on planet - removes angular momentum and drives inward migration
- Total Lindblad torque scales as:

 $\Gamma_0 = (q/h)^2 \Sigma_p r_p^4 \Omega_p^2$

- $q = M_p / M_*$ planet/star mass ratio
- Σ = gas surface density
- Ω = angular velocity
- h = aspect ratio H/r
- $X_p = X$ at the planet location.

 $\Sigma(MMSN @1AU) =>$ migration time [years] $\approx 1/q$

→ 300000 yrs for an Earth , → 20000 yrs for a Neptune !



Corotation torque

- Angular momentum is exchanged between planet and material that orbits in the horseshoe region (Goldreich & Tremaine 1980; Ward 1991, Masset 2001)
- Over one complete horseshoe orbit there is no net torque for a disc composed of ballistic particles
- Radial gradients in *entropy* and *vortensity* in a gaseous disc can give rise to a sustained corotation torque (e.g. Paardekooper et al 2010)



Corotation torque saturation - a simple argument



- We consider a disc with a negative radial entropy and temperature gradient.
- Case 1: Adiabatic evolution. The orange fluid element exchanges no heat with its surroundings no horseshoe drag
- Case 2: Locally isothermal evolution. The orange fluid element instantaneously adjusts thermally to its surroundings no horseshoe drag
- Case 3: Orange fluid element thermally equilibrates with its surroundings after 1/2 horseshoe orbit optimal corotation torque
- See Paardekooper & Mellema (2006); Baruteau & Masset (2008); Pardekooper & Papaloizou (2008); Paardekooper et al (2010, 2011)
- A similar argument applies when a vortensity gradient is present in the disc where viscosity is required for unsaturating coronation torque



Note

Viscosity and thermal diffusion are required to unsaturate the corotation torque.

Implications: Corotation torques will only be effective in disc regions where thermal or viscous diffusion operate on the appropriate time scales ~ horseshoe libration time scale

Evolution of an irradiated viscous disc model with mass ~ MMSN



Balance of Lindblad and corotation torques in irradiated viscous disc



Type II migration of high mass planets

Gap formation

Deep gap formation (δΣ/Σ < 0.1) occurs if:

 $R_{Hill} > H$ (H=disc thickness)

tidal torque > local viscous torque

 Gap formation criterion: (including pressure effects - Crida et al 2006)

q=planet-star mass ratio h= H/R (disc aspect ratio)



h = 0.05 , α_{visc} = 0.004 \rightarrow gap if q>10⁻³ .



Deep gap formation for Jupiter mass planet

Gap formation

Deep gap formation (δΣ/Σ < 0.1) occurs if:

 $R_{Hill} > H$ (H=disc thickness)

tidal torque > local viscous torque

 Gap formation criterion: (including pressure effects - Crida et al 2006)

q=planet-star mass ratio h= H/R (disc aspect ratio)

h = 0.05, $\alpha_{visc} = 0.004$

 \rightarrow gap if q>10⁻³.



Deep gap formation for Jupiter mass planet

Migration

- Type II migration occurs for a planet in a deep gap
 - migration at ~ disc viscous evolution rate (Lin & Papaloizou 1986)
- Migration rates are not precisely equal to the viscous rate (Duffel 2014; Durmann & Kley 2014)
- Large disc masses: migration rate ~ 5 x viscous rate
- Small disc masses: migration rate ~ 0.5 x viscous rate
- Detailed torque balance matters!



Gas accretion

 Simulations agree: disc supplies gas through the gap to the planet at viscous supply rate ~ 10⁻⁵ Jupiter / year (Bryden et al 1999; Kley 1999, Lubow et al 1999)

- note that numerical effects prevent accretion rate onto the planet being determined accurately! (Szulágy et al 2014)

 Gas accretion can be at a much faster rate during gap formation, building a Jovian planet in ~ 10³ yr



"Kitchen-sink" planet formation simulations

Model ingredients

- Planetary embryos + planetesimals/boulders (Mercury-6, Chambers 1996)
- Viscous disc model with stellar irradiation and photoevaporative disc wind (Lynden-Bell & Pringle 1974, Dullemond et al 2011)
- Disc cavity interior to 0.05 AU (stellar magnetosphere)
- Transition to higher disc viscosity when T > 1000 K
- Type I migration with corotation torques (Paardekooper et al 2011, Fendyke & Nelson 2014), + transition to type II migration when gap forms (Lin & Papaloizou 1986)
- Gas accretion for cores with mass > 3 Earth masses (Movshovitz et al 2010)

Model parameters

- Disc masses: 1, 1.5, 2 x MMSN
- Metallicity values: [Fe/H] = 0.5, 1, 2 x Solar
- Planetesimal/boulder radii: R_{pl} = 10m, 100m, 1km, 10km



See Hellary & Nelson (2012), Coleman & Nelson (2014), Coleman & Nelson (2016a,b for more details)

Question: Can a planet formation scenario in which planetary embryos mutually collide, accrete planetesimals/boulders and migrate through type I & II, lead to systems of planets similar to those that have been observed.

i.e. Can such a model produce the diversity of planets observed in the mass versus period diagram? Can such a model generate multiple systems of super-Earths as observed by Kepler and R.V. surveys?

Evolution in smooth discs







10

x 10⁵

Disc mass = 1 x MMSN Metallicity = 2 x solar Planetesimal sizes = 10 km



Disc mass = $1 \times MMSN$ Metallicity = $2 \times solar$ Planetesimal sizes = 1 km



Disc mass = $1 \times MMSN$ Metallicity = $2 \times solar$ Planetesimal sizes = 100 m



Disc mass = $1 \times MMSN$ Metallicity = $2 \times solar$ Planetesimal sizes = 10 m

Efficient growth and migration only occurs for small planetesimals and boulders for disc masses ~ MMSN













Forming a Jovian mass planet that orbits at \sim 5 AU requires rapid gas accretion and type II migration to initiate at \sim 14 AU

How to maintain cores at large distance and avoid rapid inward type I migration?

Evolution in radially structured discs

Radial variations in viscosity create planet traps where corotation torque prevents type I migration

Zonal flows observed in MHD simulations of disc turbulence (Papaloizou & Steinacker 2003; Johansen et al 2009; Bai & Stone 2014, Kunz & Lesur 2014; Bethune et al 2016).

Discs observed to be radially structured (e.g. ALMA partnership 2015; Zhang et al 2016)

A simple toy model: Allow viscous stress to vary by 50% at local radii - create systems of *zonal flows* with finite lifetimes



Andrews et al (2016)













Radially-structured disc

Smooth disc

Question: Can a planet formation scenario in which planetary embryos mutually collide, accrete planetesimals/boulders and migrate through type I & II, lead to systems of planets similar to those that have been observed.

i.e. Can such a model produce the diversity of planets observed in the mass versus period diagram? Can such a model generate multiple systems of super-Earths as observed by Kepler and R.V. surveys?

Answer: Yes.

But.

Radial structuring of the disc is required for giant planet formation in viscous disc models.

Resonances are much more common outcome in simulations than observed - require resonances to be unstable in the long term (see Izidoro et al 2017)